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Research Article

Littlewood-Paley g -Functions and Multipliers for the Laguerre Hypergroup

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Let $L = -(\partial^2/\partial x^2 + (2\alpha + 1/x)(\partial/\partial x) + x^2(\partial^2/\partial t^2))$; $(x, t) \in (0, +\infty) \times \mathbb{R}$, where $\alpha \geq 0$. Then L can generate a hypergroup which is called Laguerre hypergroup, and we denote this hypergroup by \mathbf{K} . In this paper, we will consider the Littlewood-Paley g -functions on \mathbf{K} and then we use it to prove the Hörmander multipliers on \mathbf{K} .

1. Introduction and Preliminaries

In [1], the authors investigated Littlewood-Paley g -functions for the Laguerre semigroup. Let

$$\mathcal{L}_\alpha = \sum_{i=1}^d x_i \frac{\partial^2}{\partial x_i^2} + (\alpha_i + 1 - x_i) \frac{\partial}{\partial x_i}, \quad (1.1)$$

where $\alpha = (\alpha_1, \dots, \alpha_d)$, $x_i > 0$, then define the following Littlewood-Paley function \mathcal{G}_α by

$$\mathcal{G}_\alpha f(x) = \left(\int_0^\infty |t \nabla_\alpha P_t^\alpha f(x)|^2 \frac{dt}{t} \right)^{1/2}, \quad (1.2)$$

where $\nabla_\alpha = (\partial_t, \sqrt{x_1} \partial_{x_1}, \dots, \sqrt{x_d} \partial_{x_d})$ and P_t^α is the Poisson semigroup associated to \mathcal{L}_α . In [1], the authors prove that \mathcal{G}_α is bounded on $L^p(\mu_\alpha)$ for $1 < p < \infty$. In this paper, we consider the following differential operator

$$L = -\left(\frac{\partial^2}{\partial x^2} + \frac{2\alpha + 1}{x} \frac{\partial}{\partial x} + x^2 \frac{\partial^2}{\partial t^2} \right); \quad (x, t) \in (0, +\infty) \times \mathbb{R}, \quad (1.3)$$

where $\alpha \geq 0$. It is well known that it can generate a hypergroup (cf. [2, 3] or [4]). We will define Littlewood-Paley g -functions associated to L and prove that they are bounded on $L^p(\mathbf{K})$ for $1 < p < \infty$. As an application, we use it to prove the Hörmander multiplier theorem on \mathbf{K} .

Let $\mathbf{K} = [0, \infty) \times \mathbf{R}$ equipped with the measure

$$dm_\alpha(x, t) = \frac{1}{\pi\Gamma(\alpha+1)} x^{2\alpha+1} dx dt, \quad \alpha \geq 0. \quad (1.4)$$

We denote by $L_\alpha^p(\mathbf{K})$ the spaces of measurable functions on \mathbf{K} such that $\|f\|_{\alpha,p} < +\infty$, where

$$\begin{aligned} \|f\|_{\alpha,p} &= \left(\int_{\mathbf{K}} |f(x, t)|^p dm_\alpha(x, t) \right)^{1/p}, \quad 1 \leq p < \infty, \\ \|f\|_{\alpha,\infty} &= \operatorname{esssup}_{(x,t) \in \mathbf{K}} |f(x, t)|. \end{aligned} \quad (1.5)$$

For $(x, t) \in \mathbf{K}$, the generalized translation operators $T_{(x,t)}^{(\alpha)}$ are defined by

$$\begin{aligned} &T_{(x,t)}^{(\alpha)} f(y, s) \\ &= \begin{cases} \frac{1}{2\pi} \int_0^{2\pi} f\left(\sqrt{x^2 + y^2 + 2xy \cos \theta}, s + t + xy \sin \theta\right) d\theta, & \text{if } \alpha = 0, \\ \frac{\alpha}{\pi} \int_0^{2\pi} \int_0^1 f\left(\sqrt{x^2 + y^2 + 2xyr \cos \theta}, s + t + xy r \sin \theta\right) r(1-r^2)^{\alpha-1} dr d\theta, & \text{if } \alpha > 0. \end{cases} \end{aligned} \quad (1.6)$$

It is known that $T_{(x,t)}^{(\alpha)}$ satisfies

$$\|T_{(x,t)}^{(\alpha)} f\|_{\alpha,p} \leq \|f\|_{\alpha,p}. \quad (1.7)$$

Let $M_b(\mathbf{K})$ denote the space of bounded Radon measures on \mathbf{K} . The convolution on $M_b(\mathbf{K})$ is defined by

$$(\mu * \nu)(f) = \int_{\mathbf{K} \times \mathbf{K}} T_{(x,t)}^{(\alpha)} f(y, s) d\mu(x, t) d\nu(y, s). \quad (1.8)$$

It is easy to see that $\mu * \nu = \nu * \mu$. If $f, g \in L_\alpha^1(\mathbf{K})$ and $\mu = f m_\alpha$, $\nu = g m_\alpha$, then $\mu * \nu = (f * g) m_\alpha$, where $f * g$ is the convolution of functions f and g defined by

$$(f * g)(x, t) = \int_{\mathbf{K}} T_{(x,t)}^{(\alpha)} f(y, s) g(y, -s) dm_\alpha(y, s). \quad (1.9)$$

The following lemma follows from (1.7).

Lemma 1.1. Let $f \in L^1_\alpha(\mathbf{K})$ and $g \in L^p_\alpha(\mathbf{K})$, $1 \leq p \leq \infty$. Then

$$\|f * g\|_{\alpha,p} \leq \|f\|_{\alpha,1} \|g\|_{\alpha,p}. \quad (1.10)$$

$(\mathbf{K}, *, i)$ is a hypergroup in the sense of Jewett (cf. [5, 6]), where i denotes the involution defined by $i(x, t) = (x, -t)$. If $\alpha = n - 1$ is a nonnegative integer, then the Laguerre hypergroup \mathbf{K} can be identified with the hypergroup of radial functions on the Heisenberg group \mathbf{H}^n .

The dilations on \mathbf{K} are defined by

$$\delta_r(x, t) = (rx, r^2t), \quad r > 0. \quad (1.11)$$

It is clear that the dilations are consistent with the structure of hypergroup. Let

$$f_r(x, t) = r^{-(2\alpha+4)} f\left(\frac{x}{r}, \frac{t}{r^2}\right). \quad (1.12)$$

Then we have

$$\|f_r\|_{\alpha,1} = \|f\|_{\alpha,1}. \quad (1.13)$$

We also introduce a homogeneous norm defined by $\|(x, t)\| = (x^4 + 4t^2)^{1/4}$ (cf. [7]). Then we can define the ball centered at $(0, 0)$ of radius r , that is, the set $B_r = \{(x, t) \in \mathbf{K} : \|(x, t)\| < r\}$.

Let $f \in L^1_\alpha(\mathbf{K})$. Set $x = \rho(\cos \theta)^{1/2}$, $t = 1/2\rho^2 \sin \theta$. We get

$$\int_{\mathbf{K}} f(x, t) dm_\alpha(x, t) = \frac{1}{2\pi\Gamma(\alpha+1)} \int_{-\pi/2}^{\pi/2} \int_0^\infty f\left(\rho(\cos \theta)^{1/2}, \frac{1}{2}\rho^2 \sin \theta\right) \rho^{2\alpha+3} (\cos \theta)^\alpha d\rho d\theta. \quad (1.14)$$

If f is radial, that is, there is a function ψ on $[0, \infty)$ such that $f(x, t) = \psi(\|(x, t)\|)$, then

$$\begin{aligned} \int_{\mathbf{K}} f(x, t) dm_\alpha(x, t) &= \frac{1}{2\pi\Gamma(\alpha+1)} \int_{-\pi/2}^{\pi/2} (\cos \theta)^\alpha d\theta \int_0^\infty \psi(\rho) \rho^{2\alpha+3} d\rho \\ &= \frac{\Gamma((\alpha+1)/2)}{2\sqrt{\pi}\Gamma(\alpha+1)\Gamma(\alpha/2+1)} \int_0^\infty \psi(\rho) \rho^{2\alpha+3} d\rho. \end{aligned} \quad (1.15)$$

Specifically,

$$m_\alpha(B_r) = \frac{\Gamma((\alpha+1)/2)}{4\sqrt{\pi}(\alpha+2)\Gamma(\alpha+1)\Gamma(\alpha/2+1)} r^{2\alpha+4}. \quad (1.16)$$

We consider the partial differential operator

$$L = -\left(\frac{\partial^2}{\partial x^2} + \frac{2\alpha+1}{x} \frac{\partial}{\partial x} + x^2 \frac{\partial^2}{\partial t^2}\right). \quad (1.17)$$

L is positive and symmetric in $L^2_\alpha(\mathbf{K})$, and is homogeneous of degree 2 with respect to the dilations defined above. When $\alpha = n - 1$, L is the radial part of the sublaplacian on the Heisenberg group \mathbf{H}^n . We call L the generalized sublaplacian.

Let $L_m^{(\alpha)}$ be the Laguerre polynomial of degree m and order α defined in terms of the generating function by

$$\sum_{m=0}^{\infty} s^m L_m^{(\alpha)}(x) = \frac{1}{(1-s)^{\alpha+1}} \exp\left(-\frac{xs}{1-s}\right). \quad (1.18)$$

For $(\lambda, m) \in \mathbf{R} \times \mathbf{N}$, we put

$$\varphi_{(\lambda, m)}(x, t) = \frac{m! \Gamma(\alpha + 1)}{\Gamma(m + \alpha + 1)} e^{i\lambda t} e^{-(1/2)|\lambda|x^2} L_m^{(\alpha)}(|\lambda|x^2). \quad (1.19)$$

The following proposition summarizes some basic properties of functions $\varphi_{(\lambda, m)}$.

Proposition 1.2. *The function $\varphi_{(\lambda, m)}$ satisfies that*

- (a) $\|\varphi_{(\lambda, m)}\|_{\alpha, \infty} = \varphi_{(\lambda, m)}(0, 0) = 1$,
- (b) $\varphi_{(\lambda, m)}(x, t) \varphi_{(\lambda, m)}(y, s) = T_{(x, t)}^{(\alpha)} \varphi_{(\lambda, m)}(y, s)$,
- (c) $L\varphi_{(\lambda, m)} = |\lambda|(4m + 2\alpha + 2)\varphi_{(\lambda, m)}$.

Let $f \in L^1_\alpha(\mathbf{K})$, the generalized Fourier transform of f is defined by

$$\widehat{f}(\lambda, m) = \int_{\mathbf{K}} f(x, t) \varphi_{(-\lambda, m)}(x, t) dm_\alpha(x, t). \quad (1.20)$$

It is easy to show that

$$\begin{aligned} (f * g)^\wedge(\lambda, m) &= \widehat{f}(\lambda, m) \widehat{g}(\lambda, m), \\ \widehat{f}_r(\lambda, m) &= \widehat{f}(r^2 \lambda, m). \end{aligned} \quad (1.21)$$

Let $d\gamma_\alpha$ be the positive measure defined on $\mathbf{R} \times \mathbf{N}$ by

$$\int_{\mathbf{R} \times \mathbf{N}} g(\lambda, m) d\gamma_\alpha(\lambda, m) = \sum_{m=0}^{\infty} \frac{\Gamma(m + \alpha + 1)}{m! \Gamma(\alpha + 1)} \int_{\mathbf{R}} g(\lambda, m) |\lambda|^{\alpha+1} d\lambda. \quad (1.22)$$

Write $L^p_\alpha(\widehat{\mathbf{K}})$ instead of $L^p(\mathbf{R} \times \mathbf{N}, d\gamma_\alpha)$. We have the following Plancherel formula:

$$\|f\|_{\alpha, 2} = \|\widehat{f}\|_{L^2_\alpha(\widehat{\mathbf{K}})}, \quad f \in L^1_\alpha(\mathbf{K}) \cap L^2_\alpha(\mathbf{K}). \quad (1.23)$$

Then the generalized Fourier transform can be extended to the tempered distributions. We also have the inverse formula of the generalized Fourier transform.

$$f(x, t) = \int_{\mathbf{R} \times \mathbf{N}} \hat{f}(\lambda, m) \varphi_{(\lambda, m)}(x, t) d\gamma_\alpha(\lambda, m) \quad (1.24)$$

provided $\hat{f} \in L^1_\alpha(\hat{\mathbf{K}})$.

In the following, we give some basic notes about the heat and Poisson kernel whose proofs can be found in [8]. Let $\{H^s\} = \{e^{-sL}\}$ be the heat semigroup generated by L . There is a unique smooth function $h((x, t), s) = h_s(x, t)$ on $\mathbf{K} \times (0, +\infty)$ such that

$$H^s f(x, t) = f * h_s(x, t). \quad (1.25)$$

We call h_s is the heat kernel associated to L . We have

$$\begin{aligned} h_s(x, t) &= \int_{\mathbf{R}} \left(\frac{\lambda}{2 \sinh(2\lambda s)} \right)^{\alpha+1} e^{-(1/2)\lambda \coth(2\lambda s)x^2} e^{i\lambda t} d\lambda, \\ h_s(x, t) &\leq C s^{-\alpha-2} e^{-(A/s)\|(x, t)\|^2}. \end{aligned} \quad (1.26)$$

Let $\{P^s\} = \{e^{-s\sqrt{L}}\}$ be the Poisson semigroup. There is a unique smooth function $p((x, t), s) = p_s(x, t)$ on $\mathbf{K} \times (0, +\infty)$, which is called the Poisson kernel, such that

$$P^s f(x, t) = f * p_s(x, t). \quad (1.27)$$

The Poisson kernel can be calculated by the subordination. In fact, we have

$$\begin{aligned} p_s(x, t) &= \frac{4s}{\sqrt{\pi}} \Gamma\left(\alpha + \frac{5}{2}\right) \int_0^\infty \left(\frac{\lambda}{\sinh \lambda} \right)^{\alpha+1} \left((s^2 + x^2 \lambda \coth \lambda)^2 + (2\lambda t)^2 \right)^{-(2\alpha+5)/4} \\ &\quad \times \cos\left(\left(\alpha + \frac{5}{2} \right) \arctan\left(\frac{2\lambda t}{s^2 + x^2 \lambda \coth \lambda} \right) \right) d\lambda, \\ p_s(x, t) &\leq C s \left(s^2 + \|(x, t)\|^2 \right)^{-(\alpha+5/2)}. \end{aligned} \quad (1.28)$$

The heat maximal function M_H is defined by

$$M_H f(x, t) = \sup_{s>0} |H^s f(x, t)| = \sup_{s>0} |(f * h_s)(x, t)|. \quad (1.29)$$

The Poisson maximal function M_P is defined by

$$M_P f(x, t) = \sup_{s>0} |P^s f(x, t)| = \sup_{s>0} |(f * p_s)(x, t)|. \quad (1.30)$$

The Hardy-Littlewood maximal function is defined by

$$M_B f(x, t) = \sup_{r>0} \frac{1}{m_\alpha(B_r)} \int_{B_r} T_{(x,t)}^{(\alpha)}(|f|)(y, s) dm_\alpha(y, s) = \sup_{r>0} (|f| * b_r)(x, t), \quad (1.31)$$

where $b(x, t) = (1/(m_\alpha(B_1)))\chi_{B_1}(x, t)$.

The following proposition is the main result of [8].

Proposition 1.3. M_B , M_P and M_B are operators on \mathbf{K} of weak type $(1, 1)$ and strong type (p, p) for $1 < p \leq \infty$.

The paper is organized as follows. In the second section, we prove that Littlewood-Paley g -functions are bounded operators on $L_\alpha^p(\mathbf{K})$. As an application, we prove the Hörmander multiplier theorem on \mathbf{K} in the last section.

Throughout the paper, we will use C to denote the positive constant, which is not necessarily same at each occurrence.

2. Littlewood-Paley g -Function on \mathbf{K}

Let $k \in \mathbb{N}$, then we define the following G -function and g_k^* -function

$$\begin{aligned} g_k(f)^2(x, t) &= \int_0^\infty \left| \partial_s^k P^s f(x, t) \right|^2 s^{2k-1} ds, \\ g_k^*(f)^2(x, t) &= \int_0^\infty \left(\int_{\mathbf{K}} s^{-(\alpha+1)} \left(1 + s^{-2} \|(y, r)\|^4 \right)^{-k} \left| \partial_s P^s T_{(y,r)}^{(\alpha)} f(x, t) \right|^2 dm_\alpha(y, r) \right) ds. \end{aligned} \quad (2.1)$$

Then, we can prove

Theorem 2.1. (a) For $k \in \mathbb{N}$ and $f \in L^2(\mathbf{K})$, there exists $C_k > 0$ such that

$$\|g_k(f)\|_{\alpha,2} = C_k \|f\|_{\alpha,2}. \quad (2.2)$$

(b) For $1 < p < \infty$ and $f \in L^p(\mathbf{K})$, there exist positive constants C_1 and C_2 , such that

$$C_1 \|f\|_{\alpha,p} \leq \|g_k(f)\|_{\alpha,p} \leq C_2 \|f\|_{\alpha,p}. \quad (2.3)$$

(c) If $k > (\alpha + 2)/2$ and $f \in L^p(\mathbf{K})$, $p > 2$, then there exists a constant $C > 0$ such that

$$\|g_k^*(f)\|_{\alpha,p} \leq C \|f\|_{\alpha,p}. \quad (2.4)$$

Proof. (a) When $k \in \mathbb{N}$, by the Plancherel theorem for the Fourier transform on \mathbf{K} ,

$$\begin{aligned}\|g_k(f)\|_{\alpha,2}^2 &= \int_{\mathbf{K}} \left(\int_0^\infty \left| \partial_s^k P^s f(x,t) \right|^2 s^{2k-1} ds \right) dm_\alpha(x,t) \\ &= \int_0^\infty \left(\int_{\mathbf{R} \times \mathbf{N}} \left| (\partial_s^k P^s f)^\sim(\lambda, m) \right|^2 d\gamma_\alpha(\lambda, m) \right) s^{2k-1} ds \\ &= \int_0^\infty \left(\int_{\mathbf{R}} \sum_{m=0}^\infty \frac{\Gamma(m+\alpha+1)}{m! \Gamma(\alpha+1)} \left| (\partial_s^k P^s f)^\sim(\lambda, m) \right|^2 |\lambda|^{\alpha+1} d\lambda \right) s^{2k-1} ds.\end{aligned}\quad (2.5)$$

Since

$$(\partial_s^k P^s f)^\sim(\lambda, m) = \hat{f}(\lambda, m) \left(-\sqrt{(4m+2\alpha+2)|\lambda|} \right)^k e^{-s\sqrt{(4m+2\alpha+2)|\lambda|}}, \quad (2.6)$$

we get

$$\begin{aligned}\|g_k(f)\|_{\alpha,2}^2 &= \int_0^\infty \left(\int_{\mathbf{R}} \sum_{m=0}^\infty \frac{\Gamma(m+\alpha+1)}{m! \Gamma(\alpha+1)} \left| \hat{f}(\lambda, m) \right|^2 ((4m+2\alpha+2)|\lambda|)^k e^{-2s\sqrt{(4m+2\alpha+2)|\lambda|}} |\lambda|^{\alpha+1} d\lambda \right) s^{2k-1} ds.\end{aligned}\quad (2.7)$$

By

$$\int_0^\infty e^{-2s\sqrt{(4m+2\alpha+2)|\lambda|}} s^{2k-1} ds = C_k ((4m+2\alpha+2)|\lambda|)^{-k}, \quad (2.8)$$

we have

$$\|g_k(f)\|_{\alpha,2}^2 = C_k \int_{\mathbf{R}} \sum_{m=0}^\infty \frac{\Gamma(m+\alpha+1)}{m! \Gamma(\alpha+1)} \left| \hat{f}(\lambda, m) \right|^2 |\lambda|^{\alpha+1} d\lambda = C_k \|f\|_{\alpha,2}^2. \quad (2.9)$$

Therefore

$$\|g_k(f)\|_{\alpha,2} = C_k \|f\|_{\alpha,2}. \quad (2.10)$$

(b) As $\{P^s\}$ is a contraction semigroup (cf. Proposition 5.1 in [3]), we can get $\|g_k(f)\|_{\alpha,p} \leq C_2 \|f\|_{\alpha,p}$ (cf. [9]). For the reverse, we can prove by polarization to the identity and (a) (cf. [10]).

(c) We first prove

$$\int_{\mathbf{K}} g_k^*(f)^2(x,t) \psi(x,t) dm_\alpha(x,t) \leq C \int_{\mathbf{K}} g_1(f)^2(x,t) M_B \psi(x,t) dm_\alpha(x,t), \quad (2.11)$$

where $0 \leq \psi \in L_\alpha^q(\mathbf{K})$ and $\|\psi\|_{\alpha,q} \leq 1$, $1/q + 2/p = 1$.

Since $k > (\alpha + 2)/2$, we know

$$\int_{\mathbf{K}} (1 + \|(y, r)\|^4)^{-k} dm_{\alpha}(y, r) < \infty. \quad (2.12)$$

By Proposition 1.3,

$$\begin{aligned} & \int_{\mathbf{K}} g_k^*(f)^2(x, t) \psi(x, t) dm_{\alpha}(x, t) \\ &= \int_{\mathbf{K}} \left(\int_0^{\infty} \int_{\mathbf{K}} s^{-(\alpha+1)} (1 + s^{-2} \|(y, r)\|^4)^{-k} \left| \partial_s P^s T_{(y,r)}^{(\alpha)} f(x, t) \right|^2 dm_{\alpha}(y, r) ds \right) \psi(x, t) dm_{\alpha}(x, t) \\ &= \int_0^{\infty} \int_{\mathbf{K}} s^{-(\alpha+1)} |\partial_s P^s f(y, r)|^2 \left(\int_{\mathbf{K}} T_{(x,t)}^{(\alpha)} (1 + s^{-2} \|(y, r)\|^4)^{-k} \psi(x, t) dm_{\alpha}(x, t) \right) dm_{\alpha}(y, r) ds \\ &\leq C \int_{\mathbf{K}} g_1(f)^2(y, r) M_B \psi(y, r) dm_{\alpha}(y, r) \\ &\leq C \|g_1(f)\|_{\alpha,p}^2 \|M_B \psi\|_{\alpha,q} \leq C \|f\|_{\alpha,p}^2. \end{aligned} \quad (2.13)$$

Therefore $\|g_k^*(f)\|_{\alpha,p} \leq C \|f\|_{\alpha,p}$. This gives the proof of Theorem 2.1. \square

We can also consider the Littlewood-Paley g -function that is defined by the heat semigroup as follows: let $k \in \mathbb{N}$, we define

$$\begin{aligned} G_k^H(f)^2(x, t) &= \int_0^{\infty} \left| \partial_s^k H^s f(x, t) \right|^2 s^{2k-1} ds, \\ G_k^{H,*}(f)^2(x, t) &= \int_0^{\infty} \left(\int_{\mathbf{K}} s^{-(\alpha+1)} (1 + s^{-2} \|(y, r)\|^4)^{-k} \left| \partial_s H^s T_{(y,r)}^{(\alpha)} f(x, t) \right|^2 dm_{\alpha}(y, r) \right) ds. \end{aligned} \quad (2.14)$$

Similar to the proof of Theorem 2.1, we can prove

Theorem 2.2. (a) For $k \in \mathbb{N}$ and $f \in L^2(\mathbf{K})$, there exists $C_k > 0$ such that

$$\|G_k^H(f)\|_{\alpha,2} = C_k \|f\|_{\alpha,2}. \quad (2.15)$$

(b) For $1 < p < \infty$ and $f \in L^p(\mathbf{K})$, there exist constants C_1 and C_2 , such that

$$C_1 \|f\|_{\alpha,p} \leq \|G_k^H(f)\|_{\alpha,p} \leq C_2 \|f\|_{\alpha,p}. \quad (2.16)$$

(c) If $k > (\alpha + 2)/2$ and $f \in L^p(\mathbf{K})$, $p > 2$, then $\|G_k^{H,*}(f)\|_{\alpha,p} \leq C \|f\|_{\alpha,p}$.

By Theorem 2.2, we can get (cf. [10])

Corollary 2.3. *Let $k \in \mathbb{N}$ and $f \in L^2(\mathbf{K})$, if $\mathcal{G}_k^H(f) \in L^p(\mathbf{K})$, $1 < p < \infty$, then $f \in L^p(\mathbf{K})$ and there exists $C > 0$ such that*

$$C\|f\|_{\alpha,p} \leq \|\mathcal{G}_k^H(f)\|_{\alpha,p}. \quad (2.17)$$

3. Hörmander Multiplier Theorem on \mathbf{K}

In this section, we prove the Hörmander multiplier theorem on \mathbf{K} . The main tool we use is the Littlewood-Paley theory that we have proved.

We first introduce some notations. Assume Ψ is a function defined on $\mathbf{R} \times \mathbf{N}$, then let $\Delta_- \Psi(\lambda, 0) = \Psi(\lambda, 0)$ and for $m \geq 1$,

$$\begin{aligned} \Delta_- \Psi(\lambda, m) &= \Psi(\lambda, m) - \Psi(\lambda, m-1), \\ \Delta_+ \Psi(\lambda, m) &= \Psi(\lambda, m+1) - \Psi(\lambda, m). \end{aligned} \quad (3.1)$$

Then we define the following differential operators:

$$\begin{aligned} \Lambda_1 \Psi(\lambda, m) &= \frac{1}{|\lambda|} (m \Delta_- \Psi(\lambda, m) + (\alpha + 1) \Delta_+ \Psi(\lambda, m)), \\ \Lambda_2 \Psi(\lambda, m) &= \frac{-1}{2\lambda} ((\alpha + m + 1) \Delta_+ \Psi(\lambda, m) + m \Delta_- \Psi(\lambda, m)). \end{aligned} \quad (3.2)$$

We have the following lemma.

Lemma 3.1. *Let $g(\lambda, m) = ((4m + 2\alpha + 2)|\lambda|)e^{-(4m+2\alpha+2)|\lambda|s}h(\lambda, m)$, where $k \in \mathbb{N}$, $h(\lambda, m)$ is a $([(\alpha + 1)/2] + 1)$ times differentiable function on \mathbb{R}^2 and satisfies*

$$\left| \left(\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) \right)^j h(\lambda, m) \right| \leq C_j ((4m + 2\alpha + 2)|\lambda|)^{-j} \quad (3.3)$$

for $j = 0, 1, 2, \dots, [(\alpha + 1)/2] + 1$. Then one has

$$\left| \left(\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) \right) g(\lambda, m) \right| \leq C \max \left\{ \frac{1}{|\lambda|^s}, 1 + \frac{m}{|\lambda|^s} \right\} e^{-\epsilon(4m+2\alpha+2)|\lambda|s}, \quad (3.4)$$

where $0 < \epsilon < 1$ and $s > 0$.

Proof. Without loss of the generality, we can assume that $\lambda > 0$. when $m = 0$, we have

$$\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) = 2 \frac{\partial}{\partial \lambda}. \quad (3.5)$$

It is easy to calculate

$$\left| \frac{\partial}{\partial \lambda} g(\lambda, 0) \right| \leq C \frac{1}{\lambda s} e^{-\epsilon(4m+2\alpha+2)\lambda s}. \quad (3.6)$$

When $m \geq 1$, we have

$$\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) = 2 \left(\frac{\partial}{\partial \lambda} - \frac{m}{\lambda} \Delta_{-1} \right). \quad (3.7)$$

Since

$$\begin{aligned} \left(\frac{\partial}{\partial \lambda} - \frac{m}{\lambda} \Delta_{-1} \right) g(\lambda, m) &= ((4m+2\alpha+2)|\lambda|) e^{-(4m+2\alpha+2)|\lambda|s} \left(\frac{\partial}{\partial \lambda} - \frac{m}{\lambda} \Delta_{-1} \right) h(\lambda, m) \\ &\quad + \frac{\partial}{\partial \lambda} \left\{ ((4m+2\alpha+2)|\lambda|) e^{-(4m+2\alpha+2)|\lambda|s} \right\} h(\lambda, m) \\ &\quad - \frac{m}{\lambda} \Delta_{-1} f(m) g(m-1), \end{aligned} \quad (3.8)$$

we get

$$\left| \left(\frac{\partial}{\partial \lambda} - \frac{m}{\lambda} \Delta_{-1} \right) g(\lambda, m) \right| \leq C \left(1 + \frac{m}{\lambda s} \right) e^{-\epsilon(4m+2\alpha+2)\lambda s}. \quad (3.9)$$

Then Lemma 3.1 is proved. \square

Then we can prove Hörmander multiplier theorem on the Laguerre hypergroup \mathbf{K} .

Theorem 3.2. *Let $h(\lambda, m)$ be a $([(\alpha+1)/2] + 1)$ times differentiable function on \mathbb{R}^2 and satisfies*

$$\left| \left(\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) \right)^j h(\lambda, m) \right| \leq C_j ((4m+2\alpha+2)|\lambda|)^{-j} \quad (3.10)$$

for $j = 0, 1, 2, \dots, [(\alpha+1)/2] + 1$ and T is an operator which is defined by $\widehat{Tf}(\lambda, m) = h(\lambda, m) \widehat{f}(\lambda, m)$, then T is bounded on $L_\alpha^p(\mathbf{K})$, where $1 < p < \infty$.

Proof. We just prove the theorem for $2 < p < \infty$, for $1 < p < 2$; we can get the result by the dual theorem. By Theorem 2.2, Corollary 2.3 and the note that $Tf \in L^2(\mathbf{K})$, it is sufficient to prove the following:

$$\mathcal{G}_2^H(Tf)(x, t) \leq C \mathcal{G}_1^{H,*}(f)(x, t), \quad (x, t) \in \mathbf{K}. \quad (3.11)$$

Let $u_s = H^s f$ and $U^s = H^s(Tf)$, then we can get

$$U^{s+t} = G_t * u_s(x, t), \quad (3.12)$$

where $\widehat{G}_t(\lambda, m) = e^{-2(2m+\alpha+1)|\lambda|t} h(\lambda, m)$.

Differentiating (3.12) with respect to t and s , then assuming that $t = s$, we can get

$$\partial_s^2 H^{2s}(Tf) = F_s * \partial_s H^s f, \quad (3.13)$$

where

$$\widehat{F}_s(\lambda, m) = -((4m + 2\alpha + 2)|\lambda|) e^{-(4m+2\alpha+2)|\lambda|s} h(\lambda, m). \quad (3.14)$$

Therefore

$$\left| \partial_s^2 H^{2s}(Tf)(x, t) \right| \leq \int_{\mathbf{K}} F_s(y, r) \left| T_{(x,t)}^{(\alpha)} \partial_s H^s f(y, r) \right| dm_\alpha(y, r). \quad (3.15)$$

By the Cauchy-Schwartz inequality,

$$\left| \partial_s^2 H^{2s}(Tf)(x, t) \right|^2 \leq A(s) \int_{\mathbf{K}} \left(1 + s^{-2} \|(y, r)\|^4 \right)^{-1} \left| T_{(x,t)}^{(\alpha)} \partial_s H^s f(y, r) \right|^2 dm_\alpha(y, r), \quad (3.16)$$

where

$$A(s) = \int_{\mathbf{K}} \left(1 + s^{-2} \|(x, t)\|^4 \right) |F_s(x, t)|^2 dm_\alpha(x, t). \quad (3.17)$$

In the following, we prove

$$A(s) \leq C s^{-\alpha-3}. \quad (3.18)$$

We write

$$\begin{aligned} A(s) &= \int_{\|(x,t)\| \leq \sqrt{s}} \left(1 + s^{-2} \|(x, t)\|^4 \right) |F_s(x, t)|^2 dm_\alpha(x, t) \\ &\quad + \int_{\|(x,t)\| > \sqrt{s}} \left(1 + s^{-2} \|(x, t)\|^4 \right) |F_s(x, t)|^2 dm_\alpha(x, t) \\ &= A_1(s) + A_2(s). \end{aligned} \quad (3.19)$$

For $A_1(s)$, we can easily get

$$\begin{aligned}
 A_1(s) &\leq C \int_{\mathbf{K}} |F_s(x, t)|^2 dm_\alpha(x, t) = C \int_{\mathbf{R} \times \mathbf{N}} \left| \widehat{F}_s(\lambda, m) \right|^2 d\gamma_\alpha(\lambda, m) \\
 &= C \int_{\mathbf{R} \times \mathbf{N}} ((4m + 2\alpha + 2)|\lambda|)^2 e^{-(8m+4\alpha+4)|\lambda|s} h^2(\lambda, m) d\gamma_\alpha(\lambda, m) \\
 &\leq C \int_{\mathbf{R}} \sum_{m=0}^{\infty} \frac{\Gamma(m + \alpha + 1)}{m! \Gamma(\alpha + 1)} ((4m + 2\alpha + 2)|\lambda|)^2 e^{-(8m+4\alpha+4)|\lambda|s} |\lambda|^{\alpha+1} d\lambda \\
 &= Cs^{-\alpha-4} \int_{\mathbf{R}} \sum_{m=0}^{\infty} \frac{\Gamma(m + \alpha + 1)}{m! \Gamma(\alpha + 1)} ((4m + 2\alpha + 2)|\lambda|)^2 e^{-(8m+4\alpha+4)|\lambda|s} |\lambda|^{\alpha+1} d\lambda \\
 &\leq Cs^{-\alpha-4} \sum_{m=0}^{\infty} (4m + 2\alpha + 2)^{-2} \leq Cs^{-\alpha-4}.
 \end{aligned} \tag{3.20}$$

For $A_2(s)$, we have

$$\begin{aligned}
 A_2(s) &\leq Cs^{-2} \int_{\mathbf{K}} (4t^2 + x^4) |F_s(x, t)|^2 dm_\alpha(x, t) \\
 &= Cs^{-2} \int_{\mathbf{K}} \left| (2it - |x|^2) F_s(x, t) \right|^2 dm_\alpha(x, t) \\
 &= Cs^{-2} \int_{\mathbf{R} \times \mathbf{N}} \left| \left(\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) \right) \widehat{F}_s(\lambda, m) \right|^2 d\gamma_\alpha(\lambda, m).
 \end{aligned} \tag{3.21}$$

By Lemma 3.1,

$$\left| \left(\Lambda_1 + 2 \left(\Lambda_2 + \frac{\partial}{\partial \lambda} \right) \right) \widehat{F}_s(\lambda, m) \right| \leq C \max \left\{ \frac{1}{|\lambda|s}, 1 + \frac{m}{|\lambda|s} \right\} e^{-\epsilon(4m+2\alpha+2)|\lambda|s}, \tag{3.22}$$

where $0 < \epsilon < 1$.

So

$$\begin{aligned}
 A_2(s) &\leq Cs^{-2} \int_{\mathbf{R} \times \mathbf{N}} e^{-\epsilon(8m+4\alpha+4)|\lambda|s} d\gamma_\alpha(\lambda, m) \\
 &= Cs^{-\alpha-4} \int_{\mathbf{R} \times \mathbf{N}} e^{-\epsilon(8m+4\alpha+4)|\lambda|s} d\gamma_\alpha(\lambda, m) \\
 &\leq Cs^{-\alpha-4}.
 \end{aligned} \tag{3.23}$$

Therefore (3.18) holds. Then

$$\left| \partial_s^2 H^{2s}(Tf)(x, t) \right|^2 \leq Cs^{-\alpha-4} \int_{\mathbf{K}} \left(1 + s^{-2} \|(y, r)\|^4 \right)^{-1} \left| T_{(x,t)}^{(\alpha)} \partial_s H^s f(y, r) \right|^2 dm_\alpha(y, r). \tag{3.24}$$

Integrating the both sides of the above inequality with $s^3 ds$, we have

$$G_2^H(x, t) \leq CG_1^{H,*}(f)(x, t). \quad (3.25)$$

Then Theorem 3.2 is proved. \square

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